## Structure and Energetics of Simple Carbenes CH<sub>2</sub>, CHF, CHCl, CHBr, CF<sub>2</sub>, and CCl<sub>2</sub>

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Abstract: There is very little experimental data concerning the singlet-triplet separations of carbenes. In many cases, it is not even possible to form estimates of this separation from experimental results.  $CH_2$  is a distinct exception; here substantial, although contradictory, data are available. The equilibrium geometries of the lowest singlet states of several carbenes have been determined experimentally. However, the geometries of many states are not known and the triplet state geometries are also generally unknown. In order to obtain theoretical information about these quantities for a series of halogenated carbenes, we have studied  $CH_2$ , CHF, CHCl, CHBr,  $CF_2$ , and  $CCl_2$  using a consistent level of ab initio electronic structure theory. One configuration (triplet state) and two configuration (singlet state) wave functions have been computed. Double  $\zeta$  basis sets augmented by a d function on the carbon atom, optimized separately for both the singlet and triplet states of each molecule, were used. The theoretical equilibrium geometries agree well with the available experimental data. The trends of the singlet-triplet states.

In recent years, one of the most interesting areas of experimental activity in organic chemistry has been carbene chemistry.<sup>2</sup> There has also been a considerable amount of theoretical research<sup>3-22</sup> concerning the simplest carbene, the methylene free radical CH<sub>2</sub>. These theoretical studies have complemented the experimental work and reliable information concerning the geometries, potential energy surfaces, and chemical reactions of the ground and excited electronic states of CH<sub>2</sub> is now available. Despite the importance of these studies, it is well to recall that in synthetic and mechanistic experimental studies substituted methylenes are often more convenient to use than the elementary CH<sub>2</sub> itself. As just one example, we point out that much of Skell's pioneering work<sup>23</sup> on the stereospecificity of carbene reactions was carried out using dibromocarbene, CBr<sub>2</sub>. However, as was the case until very recently for CH<sub>2</sub>, there is no quantitative experimental data for the singlet-triplet separation in CBr<sub>2</sub> and other halocarbenes. A knowledge of the singlet-triplet separation is essential to any reasonably complete understanding of the chemistry of a particular carbene. In the present paper, we report ab initio results for the structure and energetics of a series of halocarbenes CHX and  $CX_2$ .

Two conflicting sets of experimental values for the  ${}^{3}B_{1}-{}^{1}A_{1}$ separation,  $\Delta E$ , in CH<sub>2</sub> have been reported. The high values<sup>24-28</sup> for  $\Delta E$  are 8-9 kcal/mol (triplet state stable) and the low values<sup>29-31</sup> are 0-3 kcal/mol. Recent theoretical results<sup>18,32-35</sup> have all supported the high values. However, the very recent laser photodetachment studies of Zittel et al.<sup>36</sup> yield an even larger separation, 19.5 ± 0.7 kcal/mol. This experiment is the only "direct" measurement; i.e., it does not depend in any way upon assumed thermochemical data. Thus their result<sup>36</sup> for  $\Delta E$  could be the most accurate even though it differs substantially from all previously reported experimental values.<sup>24-31</sup>

Among the molecules studied in the present research, CHF and CF<sub>2</sub> have been the subject of two previous theoretical studies.<sup>15,37</sup> Of these, the more recent results of Staemmler<sup>15</sup> are the more reliable. Using a double  $\zeta$  plus polarization basis set, Staemmler carried out SCF and pair correlation studies of CHF and CF<sub>2</sub>. With fixed bond distances r(CH) = 1.11 Å and r(CF) = 1.30, he predicted the bond angles of the lowest singlet and triplet states, and the singlet-triplet separations, -11 (CHF) and -47 kcal (CF<sub>2</sub>), with the singlest state lying lower in both cases. These separations are not known from experiment. However, the ground state singlet structures of CHF<sup>38</sup> [r(CF) = 1.314 Å, r(CH) = 1.121 Å,  $\theta = 101.6^{\circ}$ ] and CF<sub>2</sub><sup>39-41</sup> [r(CF) = 1.300 Å,  $\theta = 104.9^{\circ}$ ] are known. The structure of the lowest singlet state of CHCl is also known from experiment,<sup>42</sup>  $r(CCl) = 1.68_9$  Å, r(CH) = 1.12 Å,  $\theta = 103.4^{\circ}$ . Although electronic spectra have been recorded for all three molecules, no triplet state transitions have been identified.

Although CCl<sub>2</sub> and CBr<sub>2</sub> have been studied extensively by matrix isolation spectroscopy,<sup>43-45</sup> their molecular structures are not known. Although the  $\Delta E$  values are also not known, a singlet ground state is presumed based on the failure to observe nonstereospecificity in the addition reactions to olefins.<sup>46</sup> Neither the vibrational nor electronic spectrum of the CHBr molecule has been observed. However, CHBr is readily generated from the corresponding diazo compound and its reactions have been investigated in some detail.<sup>2,47</sup>

**Theoretical Approach.** The present research builds upon previous theoretical work<sup>22,32-35</sup> on  $CH_2$  which suggests that quite reliable structures and reasonable singlet-triplet separations may be obtained from rather simple wave functions. For  $CH_2$ , the triplet state is approximately described by a single determinant self-consistent field (SCF) wave function.

$$1a_1^2 2a_1^2 1b_2^2 3a_1 1b_1; {}^3B_1$$
 (1)

The singlet state requires a two-configuration SCF description.

$$\begin{array}{ccc} c_1 & |a_1^2 2 a_1^2 | b_2^2 3 a_1^2 \\ \cdot e_2 & |a_1^2 2 a_1^2 | b_2^2 | b_1^2 \end{array} \ ^1A_1$$
 (2)

The key feature of this approach is that it treats the  $3a_1$  and  $1b_1$  orbitals on an equal footing for both electronic states. Near the Hartree-Fock limit (saturated basis set), this level of theory predicts<sup>22</sup> 10.9 kcal/mol for  $\Delta E$ . This value is reasonably close to that obtained from the best available configuration interaction (CI) result<sup>33</sup> ( $\Delta E = 14.1$  kcal/mol; the same authors estimate the true  $\Delta E$  value to be  $11 \pm 2$  kcal). It is also quite close to the 8-9 kcal experimental values;<sup>24-28</sup> however, it is considerably smaller than the experimental value obtained by Zittel et al.,<sup>36</sup>  $\Delta E = 19.5 \pm 0.7$  kcal/mol. In fact, their result is closer to the separation obtained when one-configuration SCF wave functions are used for both the <sup>3</sup>B<sub>1</sub> and <sup>1</sup>A<sub>1</sub> states; this computed separation<sup>22</sup> is 24.8 kcal/mol.

For all the carbenes considered in the present paper, the triplet (one configuration) and singlet (two configurations)

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Table I. Dunning's Contracted Gaussian Basis Set for the Bromine Atom<sup>a</sup>

[8s] set		[6p]	set	[2d] set		
s exponents	Coefficients	p exponents	Coefficients	d exponents	Coefficients	
439 700.	0.000 813	2957.	0.022 262	134.8	0.018 309	
66 030.	0.006 285	700.3	0.180 188	36.39	0.135 002	
15 140.	0.031 923	224.6	0.862 405	12.16	0.426 091	
4 317.	0.128 793	82.59	0.343 999	4.341	0.604 271	
1 414.	0.394 591	33.19	0.507 099	1.535	1.000 000	
523.9	0.541 292	14.20	0.258 957			
523.9	0.183 066	14.20	0.079 647			
207.7	0.617 646	7.438	0.373 442			
86.54	0.253 828	3.526	0.604 912			
30.52	1.000 000	1.595	1.000 000			
12.98	1.000 000	0.4918	1.000 000			
4.412	1.000 000	0.1507	1.000 000			
1.862	1.000 000					
0.5455	1.000 000					
0.1902	1.000 000					

<sup>a</sup> The notation used follows ref 48.

wave functions are the appropriate analogues of eq 1 and 2. (Of course CHF, CHCl, and CHBr have only a plane of symmetry and hence the corresponding term symbols are  ${}^{3}A''$  and  ${}^{1}A'$ .) The use of a two-configuration singlet and one-configuration triplet description provides a consistent picture of the singlet-triplet separation. With this relatively simple model, a study of this extended series of carbenes was computationally feasible.

The basis sets used here begin with Dunning's contracted Gaussian double  $\zeta$  sets. For hydrogen we use his (4s/2s) contraction<sup>48</sup> of Huzinaga's primitive set<sup>49</sup> (with a scale factor of 1.2); for carbon and fluorine the (9s 5p/4s 2p) contraction of Huzinaga's primitive set; for chlorine the (12s 9p/6s 4p) contraction<sup>50</sup> of Veillard's primitive set;<sup>51</sup> and for bromine Dunning's (14s 11p 5d/8s 6p 2d) contraction of his own primitive set.<sup>52</sup> Since Dunning's bromine basis has never been used or reported in the literature previously, it is given in Table I. These double  $\zeta$  bases are sufficiently flexible to allow reasonably accurate SCF descriptions of the separated atoms. The calculated, ground-state atomic SCF energies, in hartrees, with extended Slater basis near Hartree-Fock values<sup>53</sup> in parentheses, are as follows:  $H(^{2}S) = -0.4993 (-0.5); C(^{3}P) = -37.6845$ (-37.6886); F(<sup>2</sup>P) -99.3933 (-99,4093); Cl(<sup>2</sup>P) -459.4706 (-459.4819); and Br(<sup>2</sup>P) -2572.3102 (-2572.4408). With these basis sets, we varied both lengths and angles to obtain equilibrium geometries for each of the carbenes considered.

To the double  $\zeta$  basis set we added a d function centered on carbon. All six Cartesian components  $(x^2e^{-\alpha r}, xye^{-\alpha r^2},$  $xze^{-\alpha r^2}$ ,  $y^2e^{-\alpha r^2}$ , etc.) were retained in the basis set. For CH<sub>2</sub>, we simultaneously optimized both the geometry and the d function exponent  $\alpha$ . The optimum  $\alpha$  for the <sup>3</sup>B<sub>1</sub> state was 0.74; for the <sup>1</sup>A<sub>1</sub> state, it had the quite different value  $\alpha = 0.51$ . For the other carbenes, except CHBr, the singlet and triplet geometries were optimized using these CH2 d function exponents. Then, at the computed equilibrium geometries, the values of  $\alpha$  were reoptimized for each carbene. For reasons of economy, this procedure was simplified for the CHBr molecule. There the singlet and triplet geometry optimizations were carried out only at the double  $\zeta$  level. Final computations were performed at these two geometries using carbon d functions with orbital exponent  $\alpha = 0.4$ . This exponent was chosen based on the trends of optimized values for the other systems. In Table II, we give the equilibrium geometries, total energies, and mixing coefficients ( $c_1$  and  $c_2$  of eq 2), and singlet-triplet separations obtained using these basis sets.

For  $CH_2$  and  $CF_2$ , we also used more extended basis sets and optimized the geometry for these sets. The results obtained are

reported in Table II. For CH<sub>2</sub>, the addition of a set of p functions ( $\alpha = 1.0$ ) to each hydrogen changes the singlet-triplet separation from 12.3 and 12.8 kcal/mol. For both the singlet and triplet state, geometry changes are minor (bond distances change by less than 0.005 Å and bond angles by less than 1°). For difluorocarbene, the addition of a set of d functions ( $\alpha =$ 1.0) on each F atom changes  $\Delta E$  from -45.3 to -44.5 kcal. Here, however, the addition of d functions changes the bond distance more substantially. For example, the calculations including polarization functions only on carbon give  $r_e(CF)$ = 2.477  $\mu_{\rm B}$  and  $\theta_{\rm e}$  = 117.8° for the <sup>3</sup>B<sub>1</sub> state. Upon addition of polarization functions to the fluorine atoms, these change to  $r_e(CF) = 2.463 \,\mu_B$  and  $\theta_e = 118.2^\circ$  and the total energy is lowered by 0.024 hartree. Thus it is probably true that the absence of polarization functions on the terminal atoms becomes a more serious problem as one goes from  $CH_2$  to  $CF_2$ to  $CCl_2$  to  $CBr_2$ . Note also that for  $CF_2$  the carbon d exponents were reoptimized after adding the fluorine polarization functions, yielding  $\alpha = 0.69$  for the <sup>3</sup>B<sub>1</sub> state and  $\alpha = 0.67$  for the  $^{1}A_{1}$  state.

## **Results and Discussion**

We turn first to an evaluation of the molecular structures, since there is a fair amount of experimental data available for these. In Table III, we summarize the computed equilibrium geometries for the singlet and triplet states of the carbenes obtained using the double  $\zeta$  plus optimized carbon d function basis sets. Experimental values, where available, are included for comparison.

For all four experimentally known bond angles, the singlet states of CH<sub>2</sub>, CHF, CHCl, and CF<sub>2</sub>, the theoretical predictions are in close agreement with experiment, the differences being 0.1, 0.6, 1.4, and 0.6°. The best existing theoretical calculations ^33, 35 for the  ${}^3B_1$  state of  $CH_2$  give a bond angle near 134° while our SCF result is 129°. An accurate value for this bond angle is not available from experimental studies. However, matrix isolation EPR measurements<sup>54</sup> and a reinterpretation of the vacuum ultraviolet absorption spectrum<sup>55</sup> support a value<sup>54</sup> of 136  $\pm$  5°. The apparently large error of our calculated SCF bond angle is probably a result of the fact that the  ${}^{3}B_{1}$  CH<sub>2</sub> potential surface is fairly flat near equilibrium. It is interesting to note that Zittel et al.<sup>36</sup> report a still larger bond angle of  $138 \pm 4^{\circ}$ . This value is obtained by adjusting computed potential surfaces to fit the observed CH2<sup>-</sup> photodetachment spectrum.

For  $CH_2$ , CHF, and CHCl, the predicted CH bond distances agree well with experiment, the errors being 0.004,

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cule	Basis set	kcal/mol	Orb exp	<i>r</i> (CH)	r(CX)	θ	E(total)	Orb exp	R(CH)	r(CX)	θ	E(total)	с <sup>1</sup>	$c_2$
$CH_2$	No polarization functions	23.0		1.074		130.3	-38.913 69		1.099		103.5	-38.89145	0.9862	-0.165
	Optimized d function on C	12.3	0.741	1.075		128.8	-38.923 74	0.511	1.106		102.5	-38.904 11	0.9795	-0.201
	d function on C; p on H <sup>b</sup>	12.8	0.741	1.074		129.4	-38.928 22	0.511	1.103		102.8	-38.907 87		
$CF_2$	No polarization functions	-31.9			1.350	118.6	-236.595 85			1.347	103.6	-236.646 63	0.9862	-0.165
	d on carbon	-43.8	0.741		1.311	117.8	-236.648 52	0.511		1.305	104.3	-236.718 24		
	Optimized d on C	-45.3	0.65		1.311	117.8	-236.649 61	0.64		1.305	104.3	-236.721 85		
	d on C and d ( $\alpha = 1.0$ ) on F	-44.5	0.741		1.303	118.2	-236.673 26	0.511		1.291	104.7	-236.744 23	0.9842	-0.176
	Optimized d on C and d on $F^b$	b -46.5	0.69		1.303	118.2	-236.673 42	0.67		1.291	104.7	-236.747 21		
CCI3	No polarization functions	- 2.9			1.771	126.0	-956.707 34			1.830	109.5	-956.712 02	0.9852	-0.171
	d on carbon	-17.9	0.741		1.730	125.5	-956.733 10	0.511		1.756	109.4	-956.761 61		
	Optimized d on C	-13.5	0.45		1.730	125.5	-956.741 87	0.42		1.756	109.4	-956.763 48	0.9820	-0.189
CHF	No polarization functions	1.7		1.075	1.366	120.7	-137.758 22		1.107	1.373	101.9	-137.755 45	0.9863	-0.164
	d on carbon	- 7.2	0.741	1.077	1.324	120.1	-137.787 70	0.511	1.111	1.325	102.2	-137.799 20		
	Optimized d on C	- 9.2	0.66	1.077	1.321	120.4	-137.788 13	0.61	1.111	1.325	102.2	-137.802 72	0.9824	-0.186
CHCI	No polarization functions	9.2		1.075	1.779	123.8	-497.814 18		1.104	1.840	101.9	-497.799 59	0.9859	-0.167
	d on carbon	- 3.5	0.741	1.075	1.735	123.3	-497.831 91	0.511	1.101	1.762	102.0	-497.837 42		
	Optimized d on C	- 1.6	0.49	1.075	1.735	123.3	-497.835 24	0.46	1.101	1.762	102.0	-497.837 77	0.9813	-0.192
CHBr	No polarization functions	12.4		1.075	1.891	125.6	-2610.65007		1.103	1.972	102.6	-2610.632 16	0.9863	-0.164
	d on carbon	1.1	0.40	1.075	1.891	125.6	-2610.671 56	0.40	1.103	1.972	102.6	-2610.669 77	0.9812	-0.192
<i>a</i> The separati	orbital exponents of the carbon d on indicates a singlet ground state	d functions e. Bond leng	are tabulate gths are in A	d and the c	onfiguratio	n mixing co hartrees. $b_{\rm c}$	befficients, $c_1$ and $c_2 = 1.0$ .	c2 (cf. eq 2),	are given fo	or represent	ative cases	for each carbene. /	A negative si	

0.010, and 0.019 Å. As already noted for the bond angles, the errors of the calculated bond lengths become larger for the carbenes which contain heavier halogen atoms. This trend is more apparent for the CCl bond distance for the  ${}^{1}A'$  state; our ab initio value, 1.76 Å, is a full 0.07 Å longer than experiment. We suspect that the use of polarization functions (i.e., a single set of d functions) on chlorine would remove most of this error. By comparison, for the ClF<sub>2</sub> molecule, Ungemach and Schaefer<sup>56</sup> find that the addition of d functions on Cl decreases the predicted bond distances by 0.10 Å. Finally, we note that the C-F distances in singlet CHF and  $CF_2$  are predicted in very good agreement with experiment.

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For the reasons mentioned in the previous paragraph, we expect the calculated CCl distance in  $CCl_2$  to be ~0.1 Å larger than the exact (unknown) bond distance. Similarly, the calculated CBr distance in CHBr is expected to be somewhat larger than the Hartree-Fock limit (complete basis set) result. In addition, for CHBr the lack of a geometry optimization with carbon d functions also will have some effect. Comparisons for CHCl suggest that the CH distance and bond angle are relatively independent of the carbon polarization functions. However, the CCl distances in singlet and triplet CHCl were reduced by 0.08 and 0.04 Å, respectively, by geometry optimization with carbon d functions. We expect that the accuracy of the computed equilibrium geometries for the triplet states will be the same as for the singlet states discussed above.

Inspection of Table III shows several fairly clear trends. Perhaps the most obvious is the fact that the singlet bond angles are much smaller than their triplet counterparts. Less obvious but consistent with the above is the observation that with one exception  $(CF_2)$  the triplet carbene bond distances are smaller than those of the corresponding singlet states. The magnitudes of the triplet bond angles span a fairly narrow range, from 117.8 (CF<sub>2</sub>) to 128.8° (CH<sub>2</sub>). The same is true of the singlet bond angles, which vary from 102.0 (CHCl) to 109.4° (CCl<sub>2</sub>). The triplet bond angles are close to the 120° appropriate to sp<sup>2</sup> hybridization on the carbon atom. The fourth carbon electron is, of course, in the singly occupied  $p_z$  (b<sub>1</sub> or a") orbital normal to the plane of the molecule. The singlet bond angles are close to 90° which would be appropriate for a  $2s^22p_x^{-1}2p_y^{-1}$  hybridization. As we shall show later, this interpretation is supported by Mulliken population analyses of the charge on the carbon atom. Clearly, the singlet and triplet carbenes have quite different geometries.

Turning to the singlet-triplet separations, it appears clear that  $CF_2$  and  $CCl_2$  have  ${}^1A_1$  ground states. For  $CF_2$  the separation is quite large, with the singlet lying  $\sim$ 45 kcal lower. CCl<sub>2</sub> is more uncertain owing to the potential importance of Cl polarization functions but -14 kcal is a plausible estimate. For CHF, CHCl, and CHBr definitive predictions of the ground states are more difficult; the theoretical predictions are -9 kcal (singlet lower) for CHF, -2 kcal (singlet lower) for CHCl, and 1 kcal (triplet lower) for CHBr.

Between  $CH_2$  (triplet stable) and  $CF_2$  (singlet stable) there is variation of 57 kcal in the relative stabilities of the singlet and triplet states. It would be very helpful if this range of stability could be correlated to the electronegativity of the halogen member of the carbene and to the computed distribution of charge on the carbon atom. In order to provide a basis for this correlation, we give, in Table IV, Mulliken populations for the molecules studied along with calculated and estimated values of the singlet-triplet separations. The gross atomic and valence populations on carbon are given for both the singlet and triplet states. The valence population is divided into 2s (total s population -2 1s shell electrons) and 2p character. (Mulliken populations should only be used for qualitative comparisons since they tend to be basis dependent. For this reason the populations given in Table IV are those obtained with the straight double  $\zeta$  basis sets which have no polarization func-

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	T	riplet carbene			Singlet carbene	
Molecule	<i>r</i> (CH)	r(CX)	θ	r(CH)	<i>r</i> (CX)	θ
CH <sub>2</sub>	1.075		128.8	1.106 (1.11 <sup>a</sup> )		102.5 (102.4 <i><sup>a</sup></i> )
CHF	1.077	1.321	120.4	$1.111(1.121^{b})$	$1.325(1.314^{\circ})$	$102.2(101.6^{\circ})$
CHCl	1.075	1.735	123.3	$1.101(1.12^d)$	$1.762(1.689^d)$	$102.0(103.4^{d})$
CHBr	1.075	1.891	125.6	1.103	1.972	102.6
CF <sub>2</sub>		1.311	117.8		1.305 (1.300 <sup>e</sup> )	104.3 (104.9 <sup>e</sup> )
CCl <sub>2</sub>		1.730	125.5		1.756	109.4

<sup>a</sup> G. Herzberg, "Electronic Spectra of Polyatomic Molecules", Van Nostrand, Princeton, N.J., 1967. <sup>b</sup> Assumed. <sup>c</sup> Reference 38. <sup>d</sup> Reference 42. <sup>e</sup> Reference 41. <sup>f</sup> Bond distances are reported in Å and bond angles in degrees. The results in this table were obtained using double  $\zeta$  basis sets augmented by an optimized d function on the carbon atom. Experimental values are given in parentheses.

Table IV. Mulliken Populations and Singlet-Triplet Separations  $\Delta E$  for Several Carbenes<sup>a</sup>

	Triplet population		Sing	Singlet population				Calcd $\Delta E$ ,	
	Total	2s	2p	Total	2s	2p	<i>R</i> <sub>1</sub>	<i>R</i> <sub>2</sub>	kcal/mol
CH <sub>2</sub>	6.36	1.43	2.93	6.25	1.71	2.53	1.20	1.23	+12.8
CHBr	6.25	1.48	2.77	6.13	1.83	2.30	1.24	1.27	+1.1
CHCl	6.14	1.45	2.69	6.03	1.81	2.21	1.25	1.29	-1.6
CHF	5.85	1.33	2.53	5.77	1.72	2.05	1.29	1.32	-9.2
CCl <sub>2</sub>	5.99	1.45	2.53	5.89	1.87	2.02	1.29	1.32	-13.5
$CF_2$	5.39	1.12	2.28	5.41	1.69	1.71	1.51	1.51	-44.5

<sup>a</sup> The calculated values of  $\Delta E$  are those obtained with the basis set which includes optimized carbon d functions. The quantities  $R_1$  and  $R_2$  are ratios of 2s and 2p populations for the singlet and triplet states and are defined in eq 3 and 4.

tions included.) The total carbon populations are nearly the same, within 0.1 electron, for the singlet and triplet states of each carbene. The division into 2s and 2p character is, however, quite different. For the triplet states, the 2s population varies from 1.12 (CF<sub>2</sub>) to 1.43 (CH<sub>2</sub>) suggesting a carbon hybridization approaching sp<sup>2</sup> (with an additional singly occupied  $p_z$  orbital) and consistent with the computed bond angles near 120° (cf. Table III). For the singlet states, the 2s and 2p populations are more nearly equal and usually close to 2 consistent with a hybridization approaching s<sup>2</sup>p<sub>x</sub>p<sub>y</sub> and nearly 90° bond angles.

The singlet-triplet separation correlates with the total charge on carbon. When carbon has a negative charge, the carbene is triplet stable; when it has a positive charge, the carbene is singlet stable. Both chemical intuition and the observed Mulliken populations suggest that fluorine is more electronegative than chlorine, and chlorine more so than bromine. The singlet-triplet separation is greatest for CF<sub>2</sub>, and CCl<sub>2</sub> is also predicted to have a ground-state singlet. Although we did not perform comparable studies of CBr<sub>2</sub>, it seems likely that the <sup>1</sup>A<sub>1</sub> and <sup>3</sup>B<sub>1</sub> states will be nearly degenerate, lying within a few kcal of each other. In light of the greater reactivity<sup>14,20,21</sup> of singlet carbenes, it is not surprising that CBr<sub>2</sub> appears to have the characteristics of a ground-state singlet.<sup>23</sup> The CHX species exhibit the same general trend.

The dependence of the singlet-triplet separation on the charge on carbon can be understood in terms of the very different hybridization for the singlet and triplet states. In order to form an sp<sup>2</sup> hybrid a rather large promotion energy is required; the energy difference<sup>57</sup> between  $C(2s^2 2p^2; ^3P)$  and  $C(2s2p^3; ^5S)$  is 4.2 eV. This promotion energy is compensated by the strength of the bonds formed with the sp<sup>2</sup> hybrids. When the carbon is negatively charged, as in the case for CH<sub>2</sub>, CHBr, and CHCl, some of the charge goes into the atomic 2s orbital reducing the "promotion energy". When electrons are removed from carbon by strongly electronegative neighbors (the case for CHF, CCl<sub>2</sub>, and CF<sub>2</sub>), the "promotion energy" is increased. For example, the energy difference<sup>57</sup> between C<sup>+</sup> (2s<sup>2</sup>2p;<sup>2</sup>P) and C<sup>+</sup> (2s<sup>2</sup>2p;<sup>2</sup>P) is 5.3 eV. In this case, the triplet becomes less stable with respect to the singlet state. For the singlet

states, the 2s population is much closer to 2 and the promotion energies are correspondingly smaller than for the triplets. Of course, the bonds formed are also relatively weaker. We can form two ratios to measure the relative values of the promotion energies for the singlet and triplet states:

$$R_1 = \frac{\text{singlet 2s population}}{\text{triplet 2s population}}$$
(3)

and

$$R_2 = \frac{\text{singlet } 2\text{s pop/singlet } (2\text{s} + 2\text{p}) \text{ pop}}{\text{triplet } 2\text{s pop/triplet } (2\text{s} + 2\text{p}) \text{ pop}}$$
(4)

The smaller the ratio, the smaller the promotion energy difference will be between the triplet and singlet states. Hence, the triplet state will be more stable than the singlet. The reverse will be true for larger ratios. The values of  $R_1$  and  $R_2$  are given in Table IV and the expected behavior is, indeed, followed.

Goddard et al.<sup>58</sup> have considered qualitatively the relative stability of the singlet and triplet states for CH<sub>2</sub> and CF<sub>2</sub> using a valence bond analysis. Their arguments are directly related to the excitation energies of the <sup>4</sup>Σ<sup>-</sup> STATES OF CH and CF. In order to form the triplet state of CX<sub>2</sub> (X = H or F), they find that one C-X bond must be like that in the <sup>4</sup>Σ<sup>-</sup> state of the corresponding diatomic molecule. The other C-X bond is like that in the ground <sup>2</sup>II state. For the singlet state of CX<sub>2</sub>, both C-X bonds are like that in the <sup>2</sup>II state. The <sup>4</sup>Σ<sup>-</sup> state of CH is close to, ~0.5 eV above, the <sup>2</sup>II state; <sup>59,60</sup> the <sup>4</sup>Σ<sup>-</sup> state of CH is highly excited, <sup>58</sup> ~3 eV above <sup>2</sup>II. Thus the triplet state of CF<sub>2</sub>, which requires a <sup>4</sup>Σ<sup>-</sup>-like bond, is expected to be considerably higher than the singlet state. This is found to be the case in our calculations and in previous work.<sup>15,37</sup>

However, a simple extension of the arguments of Goddard et al.<sup>55</sup> does not explain the different singlet-triplet separations found for CH<sub>2</sub> and CHF. In both the singlet and triplet states of CHF, the C-F bond would be the favorable <sup>2</sup>II-like bond; only the C-H bond would change from <sup>4</sup>Σ<sup>-</sup>-like (triplet CHF) to <sup>2</sup>II-like (singlet CHF). Thus CH<sub>2</sub> and CHF would be expected to have similar singlet-triplet separations. In this and previous work,<sup>15,37</sup> however, these separations are found to be quite different; in every case,  $\Delta E_{ST}(CH_2) - \Delta E_{ST}(CHF)$  is

found to be  $\sim 22$  kcal/mol. Our correlation of the singlettriplet separation with the Mulliken atomic populations for C (and hence the electronegativity of the halogen substituent) is able to account for the separation in CHF as well as in CF<sub>2</sub>.

Finally, we note that all of the above halocarbenes (plus CBr<sub>2</sub>) have been studied by Hoffman, Zeiss, and Van Dine<sup>61</sup> using extended Hückel theory. These authors emphasize that their results should be taken qualitatively. In fact, many of the trends reported here are seen in this earlier semiempirical study.

Note Added in Proof. We have recently (unpublished) carried out very large scale CI studies of the CH<sub>2</sub> singlet-triplet separation. It now seems clear that the  $\Delta E$  value of  $11 \pm 2$  kcal predicted in ref 33 is correct. This conclusion greatly strengthens the credibility of the much simpler molel studies reported in the present paper.

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